



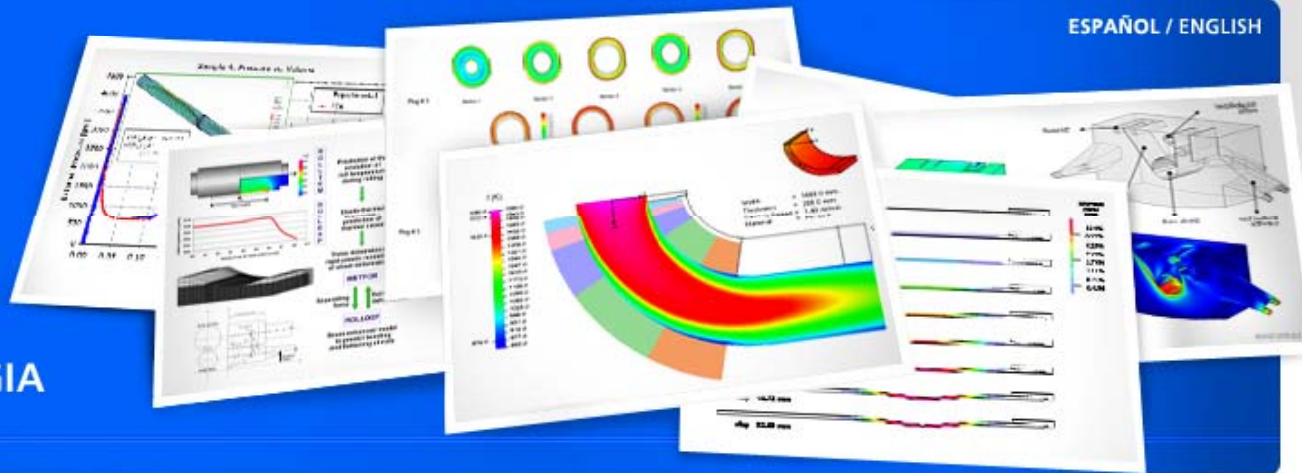
# SIM&TEC

Simulación y Tecnología

*Simulation and Technology*

ESPAÑOL / ENGLISH

DE LA CIENCIA  
A LA TECNOLOGIA



# Advanced Topics in Computational Solid Mechanics.

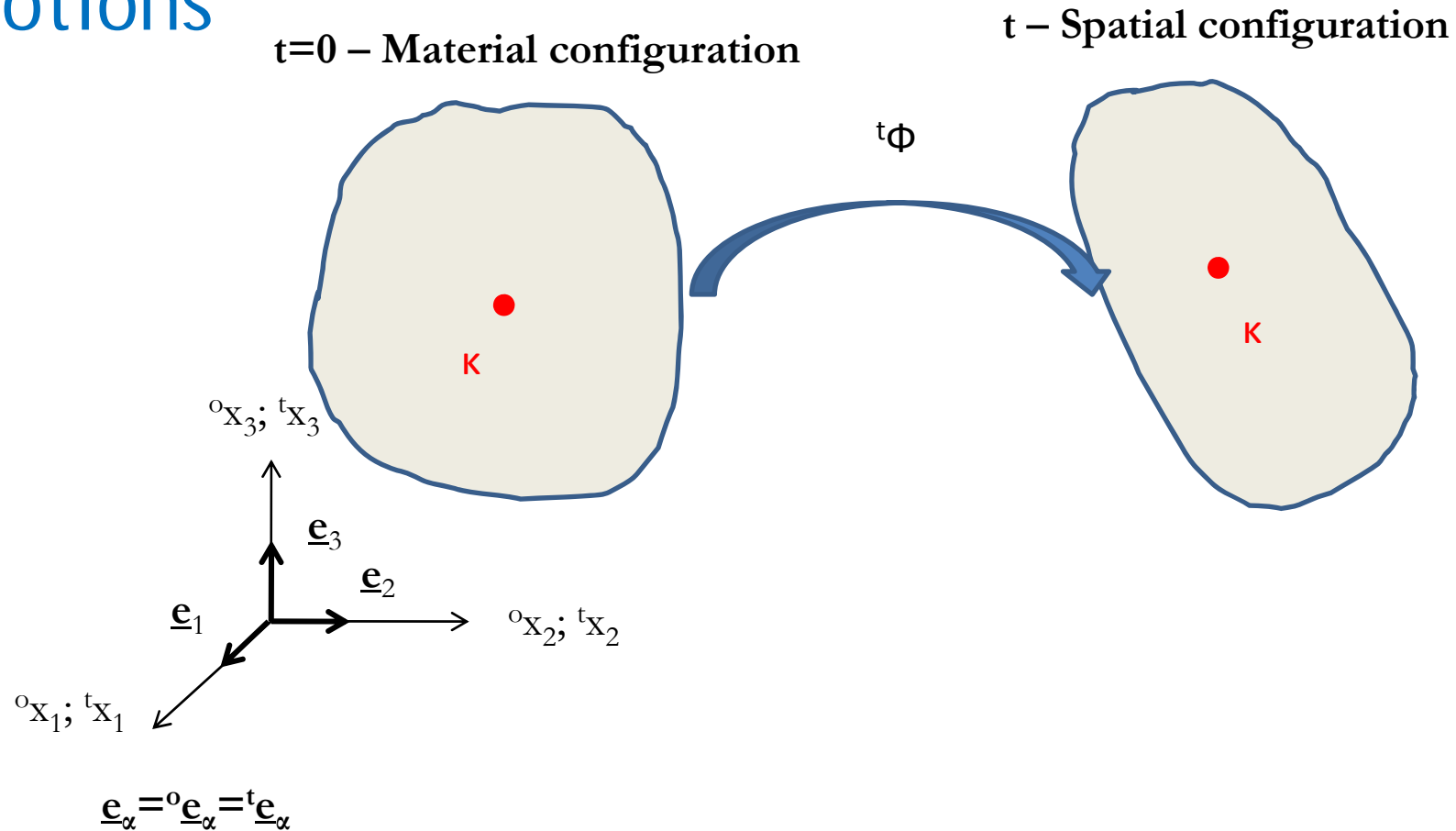
## Section 2: Kinematics of the continuous media

Eduardo N. Dvorkin

Stanford University  
Mechanical Engineering  
Winter Quarter 2010

# The continuum media and its configuration.

## Motions



# The continuum media and its configuration

## Motions

$$\begin{array}{l}
 {}^t x_a = {}^t x_a(\chi, t) \quad ; \quad \chi = \chi({}^t x_a) \\
 {}^t \tilde{x}_b = {}^t \tilde{x}_b({}^t x_a) \quad ; \quad {}^t x_a = {}^t x_a({}^t \tilde{x}_b) \\
 {}^o x_A = {}^o x_A(\chi) \quad ; \quad \chi = \chi({}^o x_A)
 \end{array}
 \left. \begin{array}{l} \\ \\ \end{array} \right\} \begin{array}{l} \text{Spatial configuration} \\ \\ \text{Material configuration} \\ \text{Reference configuration} \end{array}$$

Motion:

$$\begin{array}{l}
 {}^t x_a = {}^t \phi_a({}^o x_A, t) \\
 {}^t \tilde{x}_a = {}^t \tilde{x}_a({}^t x_b) = {}^t \tilde{x}_a[{}^t \phi_b({}^o x_A)] = {}^t \tilde{\phi}_a({}^o x_A)
 \end{array}$$

Regular motion : no opening of holes and no material interpenetration

$${}^o x_A = [{}^t \phi^{-1}]_A({}^t x_A)$$

---

# The continuum media and its configuration Motions

$${}^o \underline{x}(\boldsymbol{\kappa}) = {}^o x_a(\boldsymbol{\kappa}) {}^o \underline{e}_a \quad ; \quad {}^t \underline{x}(\boldsymbol{\kappa}) = {}^t x_a(\boldsymbol{\kappa}) {}^t \underline{e}_a$$

$${}^t \underline{u}_a = {}^t x_a - {}^o x_a$$

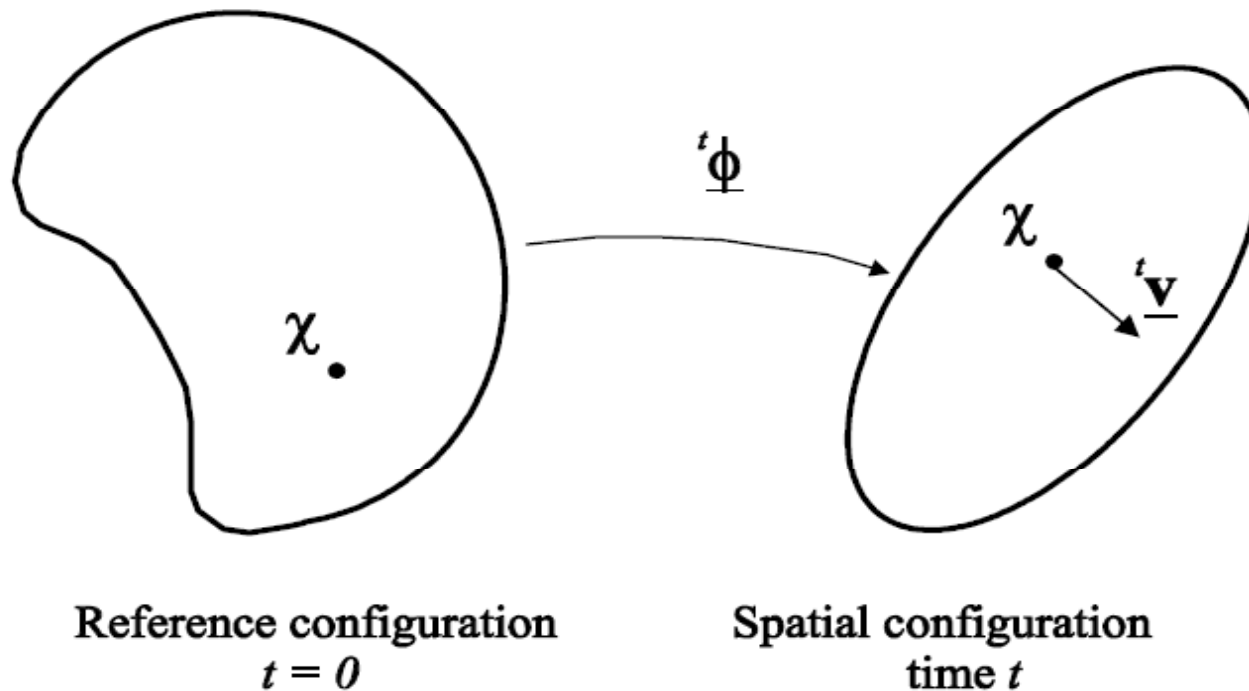
Material velocity:

$${}^t \underline{\mathbf{v}}(\boldsymbol{\chi}, t) = \frac{\partial {}^t \underline{\mathbf{x}}(\boldsymbol{\chi}, t)}{\partial t}$$

# The continuum media and its configuration Motions

<i>Lagrangean (material) description of motion</i>	${}^t \underline{v} = {}^t \underline{v}({}^o \mathbf{x}_A, t)$
<i>Eulerian (spatial) description of motion</i>	${}^t \underline{v} = {}^t \underline{v}({}^t \mathbf{x}_a, t)$

# The continuum media and its configuration Motions



Material velocity of a particle

$${}^t \underline{v} = {}^t v_\alpha \mathbf{e}_\alpha$$

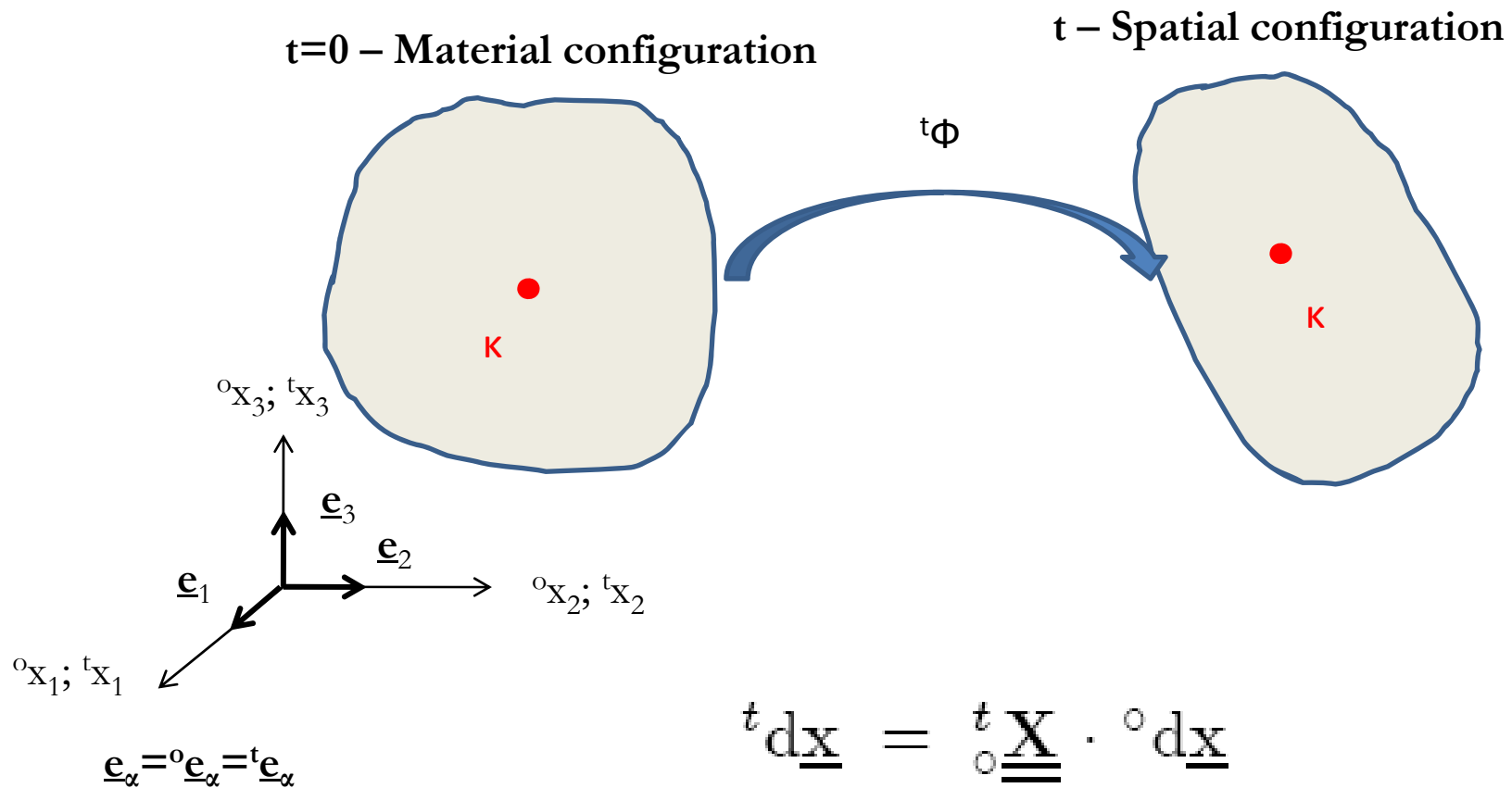
# The continuum media and its configuration Motions

Temporal material derivative: following the particle  ${}^t \underline{a} = \frac{D^t \underline{v}}{Dt}$ .

<i>Description</i>	<i>Coordinates</i>	
<i>Lagrangean (material) description of motion</i>	<i>Fixed Cartesian</i>	${}^t \underline{a} = \frac{\partial^t v_\alpha}{\partial t} \underline{e}_\alpha$
<i>Eulerian (spatial) description of motion</i>	<i>Fixed Cartesian</i>	${}^t \underline{a} = \left[ \frac{\partial^t v_\alpha}{\partial t} + \frac{\partial^t v_\alpha}{\partial^t z_\beta} {}^t v_\beta \right] \underline{e}_\alpha$



# The deformation gradient tensor



# The deformation gradient tensor

$${}^t\underline{\underline{X}} = \frac{\partial {}^t x_a}{\partial {}^o x_A} {}^t \underline{e}_a {}^o \underline{e}_A$$

The deformation gradient tensor is a two-point tensor

$${}^o d\underline{x} = {}^t\underline{\underline{X}}^{-1} \cdot {}^t d\underline{x}$$

$${}^t\underline{\underline{X}}^{-1} = \frac{\partial {}^o x_A}{\partial {}^t x_a} {}^o \underline{e}_A {}^t \underline{e}_a$$

# The deformation gradient tensor

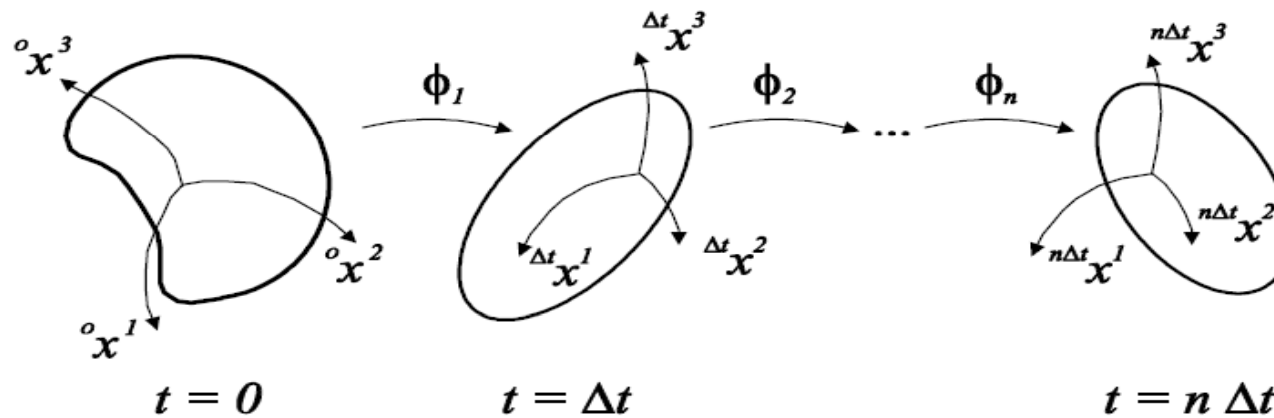


Fig. 2.3. Sequence of motions

$${}_{0}^{n\Delta t}X_{aP} = \frac{\partial^{n\Delta t}x_a}{\partial {}^0x_P} = \frac{\partial^{n\Delta t}x_a}{\partial^{(n-1)\Delta t}x_b} \cdots \frac{\partial^{\Delta t}x_l}{\partial {}^0x_P}$$

Therefore,

$${}_{0}^{n\Delta t}\underline{\underline{X}} = {}_{(n-1)\Delta t}^{n\Delta t}\underline{\underline{X}} \cdot {}_{(n-2)\Delta t}^{(n-1)\Delta t}\underline{\underline{X}} \cdots {}_{\Delta t}^{2\Delta t}\underline{\underline{X}} \cdot {}_0^{\Delta t}\underline{\underline{X}}.$$

---

# The deformation gradient tensor

$${}^t J(\chi, t) = \frac{{}^t dV}{{}^o dV}$$

$${}^t J(\chi, t) = |{}^t X|$$

# The right polar decomposition

The Green deformation tensor

$${}^t\underline{\underline{\mathbf{C}}} = {}^t\underline{\underline{\mathbf{X}}}^T \cdot {}^t\underline{\underline{\mathbf{X}}}$$

$${}^t d\underline{\underline{\mathbf{x}}} \cdot {}^t d\underline{\underline{\mathbf{x}}} = {}^o d\underline{\underline{\mathbf{x}}} \cdot {}^t\underline{\underline{\mathbf{C}}} \cdot {}^o d\underline{\underline{\mathbf{x}}}$$

- Symmetric
- Positive-definite

# The right polar decomposition

The right stretch tensor

$${}^t\underline{\underline{\mathbf{U}}} = [{}^t\underline{\underline{\mathbf{C}}}]^{1/2}$$

- Symmetric
- Positive-definite

The right polar decomposition

$${}^t\underline{\underline{\mathbf{X}}} = {}^t\underline{\underline{\mathbf{R}}} \cdot {}^t\underline{\underline{\mathbf{U}}}$$

$$(i) \quad {}^t\underline{\underline{\mathbf{R}}}^T \cdot {}^t\underline{\underline{\mathbf{R}}} = \underline{\underline{\mathbf{1}}}$$

$$(ii) \quad {}^t\underline{\underline{\mathbf{R}}} \cdot {}^t\underline{\underline{\mathbf{R}}}^T = \underline{\underline{\mathbf{1}}}$$

${}^t\underline{\underline{\mathbf{R}}}$  is orthogonal

The right polar decomposition is unique

# The right polar decomposition

## Physical interpretation

$$\text{If } {}^t\underline{\underline{\mathbf{X}}} = {}^t\underline{\underline{\mathbf{R}}}$$

$${}^t d\underline{\underline{\mathbf{x}}}_1 \cdot {}^t d\underline{\underline{\mathbf{x}}}_2 = {}^o d\underline{\underline{\mathbf{x}}}_1 \cdot {}^o d\underline{\underline{\mathbf{x}}}_2$$

$${}^t d\underline{\underline{\mathbf{x}}}_1 \cdot {}^t d\underline{\underline{\mathbf{x}}}_2 = {}^o d\underline{\underline{\mathbf{x}}}_1 \cdot {}^t\underline{\underline{\mathbf{C}}} \cdot {}^o d\underline{\underline{\mathbf{x}}}_2 = {}^o d\underline{\underline{\mathbf{x}}}_1 \cdot ({}^t\underline{\underline{\mathbf{U}}} \cdot {}^t\underline{\underline{\mathbf{U}}}) \cdot {}^o d\underline{\underline{\mathbf{x}}}_2$$

---

# The left polar decomposition

The Finger deformation tensor

$$\underline{\underline{t}}{\mathbf{b}} = \underline{\underline{t}}{\mathbf{X}} \cdot \underline{\underline{t}}{\mathbf{X}}^T$$

- Symmetric
- Positive-definite



---

# The left polar decomposition

$${}^t\underline{\underline{\mathbf{X}}} = {}^t\underline{\underline{\mathbf{V}}} \cdot {}^t\underline{\underline{\mathbf{R}}}$$

The left stretch tensor

$${}^t\underline{\underline{\mathbf{V}}} = [{}^t\underline{\underline{\mathbf{b}}}]^{1/2}$$

The left polar decomposition is unique

---

# The left polar decomposition

## Physical interpretation

$${}^o d\underline{\mathbf{x}}_1 \cdot {}^o d\underline{\mathbf{x}}_2 = {}^t d\underline{\mathbf{x}}_1 \cdot \left( {}^t \underline{\underline{\mathbf{V}}}^{-1} \cdot {}^t \underline{\underline{\mathbf{V}}}^{-1} \right) \cdot {}^t d\underline{\mathbf{x}}_2$$

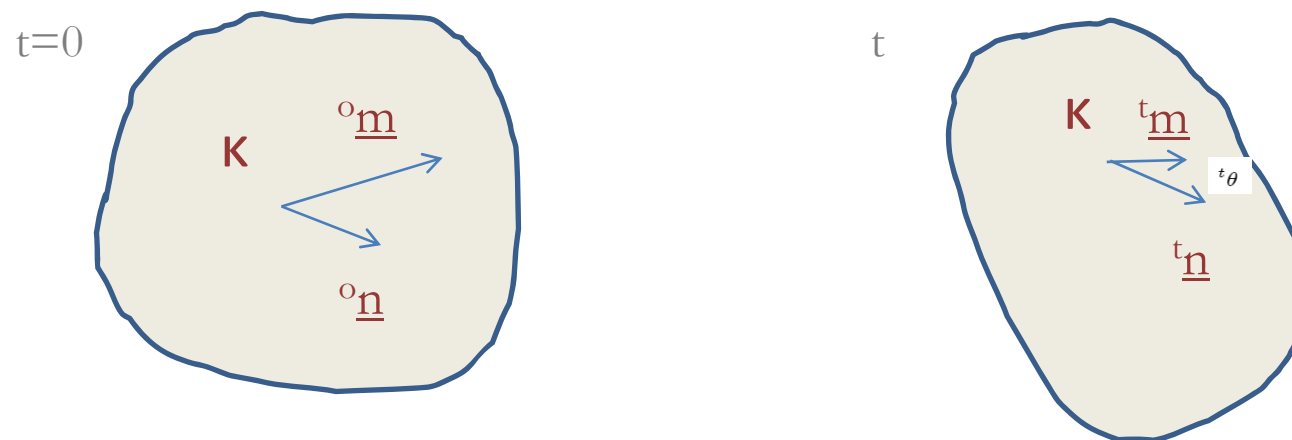
Please notice that:

$${}^t \underline{\underline{\mathbf{V}}}^T = {}^t \underline{\underline{\mathbf{R}}} \cdot {}^t \underline{\underline{\mathbf{U}}} \cdot {}^t \underline{\underline{\mathbf{R}}}^T = {}^t \underline{\underline{\mathbf{V}}}$$

${}^t \underline{\underline{\mathbf{U}}}$  and  ${}^t \underline{\underline{\mathbf{V}}}$  have the same eigenvalues

# The polar decomposition

## Physical interpretation



$$\frac{{}^t dS}{{}^o dS} = \left[ [{}^o n]^T [{}^t C] [{}^o n] \right]^{1/2}$$

$${}^t \theta = \cos^{-1} \left[ \frac{[{}^o m]^T [{}^t C] [{}^o n]}{\left( \frac{{}^t dS}{{}^o dS} \right)_m \left( \frac{{}^t dS}{{}^o dS} \right)_n} \right]$$

# Strain measures

*Green*  ${}^t d\underline{\mathbf{x}}_1 \cdot {}^t d\underline{\mathbf{x}}_2 = {}^o d\underline{\mathbf{x}}_1 \cdot {}^t \underline{\underline{\mathbf{C}}} \cdot {}^o d\underline{\mathbf{x}}_2$

*Finger*  ${}^t d\underline{\mathbf{x}}_1 \cdot {}^t \underline{\underline{\mathbf{b}}}^{-1} \cdot {}^t d\underline{\mathbf{x}}_2 = {}^o d\underline{\mathbf{x}}_1 \cdot {}^o d\underline{\mathbf{x}}_2$

$${}^t d\underline{\mathbf{x}}_1 \cdot {}^t d\underline{\mathbf{x}}_2 - {}^o d\underline{\mathbf{x}}_1 \cdot {}^o d\underline{\mathbf{x}}_2 = 2 {}^o d\underline{\mathbf{x}}_1 \cdot \underbrace{\left[ \frac{1}{2} ({}^t \underline{\underline{\mathbf{C}}} - {}^o \underline{\underline{\mathbf{g}}}) \right]}_{{}^t \underline{\underline{\epsilon}}}$$

*Green-Lagrange*  ${}^t \underline{\underline{\epsilon}}$

# Strain measures

$${}^t d\underline{\mathbf{x}}_1 \cdot {}^t d\underline{\mathbf{x}}_2 - {}^o d\underline{\mathbf{x}}_1 \cdot {}^o d\underline{\mathbf{x}}_2 = 2 \, {}^t d\underline{\mathbf{x}}_1 \cdot \underbrace{\left[ \frac{1}{2} ({}^t \underline{\mathbf{g}} - {}^t \underline{\mathbf{b}}^{-1}) \right]}_{{}^t \underline{\mathbf{e}}} \cdot {}^t d\underline{\mathbf{x}}_2$$

*Almansi*  ${}^t \underline{\mathbf{e}}$

*Hencky*  ${}^t \underline{\mathbf{H}} = \ln {}^t \underline{\mathbf{U}}$

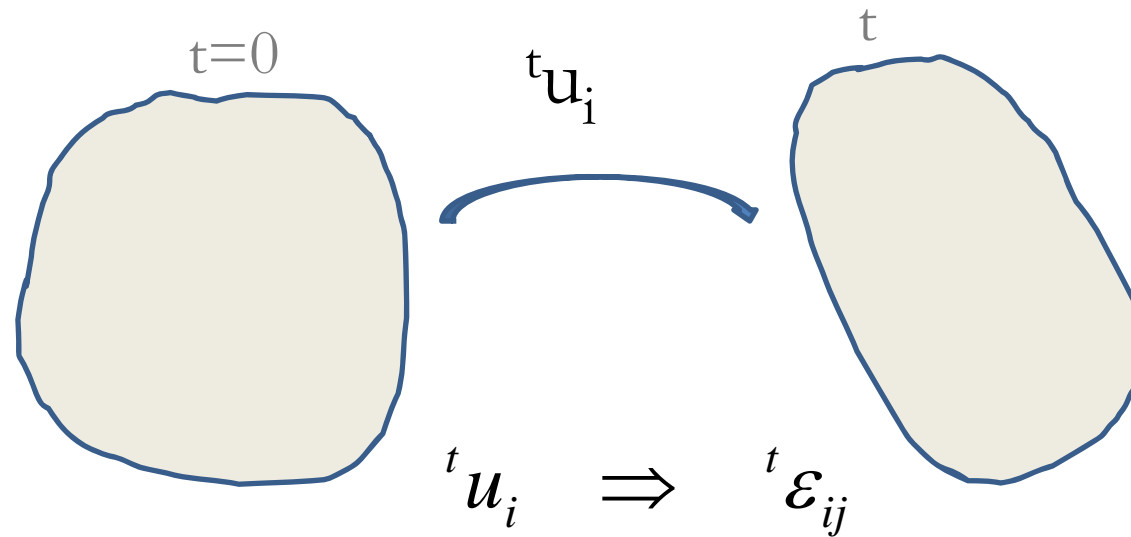
# Strain rates

*Velocity gradient tensor* 
$${}^t \underline{l} = \frac{\partial^t v_i}{\partial^t x_i} {}^t \underline{e}_i {}^t \underline{e}_i$$

$${}^t \underline{\underline{l}} = {}^t \underline{\underline{d}} + {}^t \underline{\underline{\omega}}$$

Strain rate tensor	${}^t \underline{\underline{d}} = {}^t \underline{\underline{d}}^T = \frac{1}{2} ({}^t \underline{\underline{l}} + {}^t \underline{\underline{l}}^T)$
Vorticity tensor	${}^t \underline{\underline{\omega}} = - {}^t \underline{\underline{\omega}}^T = \frac{1}{2} ({}^t \underline{\underline{l}} - {}^t \underline{\underline{l}}^T)$

# Compatibility



${}^t \epsilon_{ij} \Rightarrow {}^t u_i$  If

$$\frac{\partial^2 {}^t \epsilon_{\alpha\beta}}{\partial^{\circ} z^{\gamma} \partial^{\circ} z^{\delta}} + \frac{\partial^2 {}^t \epsilon_{\gamma\delta}}{\partial^{\circ} z^{\alpha} \partial^{\circ} z^{\beta}} - \frac{\partial^2 {}^t \epsilon_{\alpha\delta}}{\partial^{\circ} z^{\gamma} \partial^{\circ} z^{\beta}} - \frac{\partial^2 {}^t \epsilon_{\gamma\beta}}{\partial^{\circ} z^{\alpha} \partial^{\circ} z^{\delta}} = 0$$